

# The Coriolis Force

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## Introduction

Imagine you're on a rotating merry-go-round, and you throw a ball to another person who's on the opposite side of the merry-go-round. If you aim directly at the other person, you'll miss them—the ball will travel in a straight line relative to the ground, but the merry-go-round will have rotated during the time the ball is in the air. Relative to the merry-go-round, the ball will appear to move along a curved path. You can attribute this curvature to a “fictitious force” called the *Coriolis force*. The Coriolis force is not a real force—it's just an artifact of viewing the ball's motion in a rotating reference frame. The ball really moves in a straight line relative to the ground.

But suppose you're in the rotating reference frame of the merry-go-round. You'll see the ball move in a curved path, which can't happen unless there is a “force” present. We can compute the magnitude of this Coriolis force by considering the following situation. Suppose you're at the center of the merry-go-round, and throw a ball outward with velocity  $v$  while the merry-go-round is rotating with an angular velocity  $\Omega$ . After a time  $t$ , the ball will have moved a radial distance  $r = vt$ . At time  $t$ , a point on the merry-go-round a distance  $r$  from the center will have moved a distance of arc length

$$s = r\theta \tag{1}$$

$$= r(\Omega t) \tag{2}$$

$$= (vt)\Omega t \tag{3}$$

$$= \Omega vt^2. \tag{4}$$

But under a constant acceleration  $a_c$ , we know

$$s = \frac{1}{2}a_c t^2. \tag{5}$$

Comparing Eq. (4) with Eq. (5), we deduce that the Coriolis acceleration  $a_c$  is given by

$$a_c = 2\Omega v. \tag{6}$$

More generally, in terms of vectors, the Coriolis acceleration vector  $\mathbf{a}_c$  is given by

$$\boxed{\mathbf{a}_c = -2(\boldsymbol{\Omega} \times \mathbf{v})} \tag{7}$$

From Newton's Second Law, the corresponding Coriolis force  $\mathbf{F}_c$  on a body of mass  $m$  is then

$$\boxed{\mathbf{F}_c = -2m(\boldsymbol{\Omega} \times \mathbf{v})} \tag{8}$$

## Examples

### Golf

For example, suppose we're on the surface of the Earth, in the northern hemisphere, and hit a golf ball due south with velocity  $v$ . Since the Earth rotates to the east, the Earth's angular velocity vector  $\Omega$  is along the Earth's axis, northward out of the north pole. Then by Eq. (8), there will be a *westward* Coriolis force acting on the golf ball, equal in magnitude to

$$F_c = 2 m \Omega v \sin \varphi, \quad (9)$$

where  $\varphi$  is the latitude and  $m$  is the mass of the golf ball. This will cause the ball to slice the right. The effect is very slight, though. For example, given the rotation rate of the Earth  $\Omega = 7.2921 \times 10^{-5}$  rad/s, the mass of the golf ball  $m = 45$  g, a typical ball speed  $v = 50$  m/s, and a latitude of  $\varphi = 39^\circ$ , the Coriolis force only amounts to  $F_c = 206.5 \mu\text{N}$ , or about 0.05% of the weight of the golf ball.

The Coriolis force is zero at the equator, and greater at higher latitudes. In the southern hemisphere, the Coriolis force will cause a slight hook of the ball to the left, rather than the slice it will experience in the northern hemisphere.

### Weather

By Eq. (8), we can see that in the northern hemisphere, air currents moving northward are deflected to the east; eastward currents are deflected to the south; southward currents are deflected to the west; and westward currents are deflected to the north. If a low-pressure area forms in the atmosphere, then the pressure gradients will cause the air currents to flow toward the center of the area; but because of the Coriolis deflections, the result will be that the air currents will flow counter-clockwise, creating an air pattern called a *cyclone* around the low-pressure area. Similarly, in the southern hemisphere, cyclones will be air currents rotating clockwise.